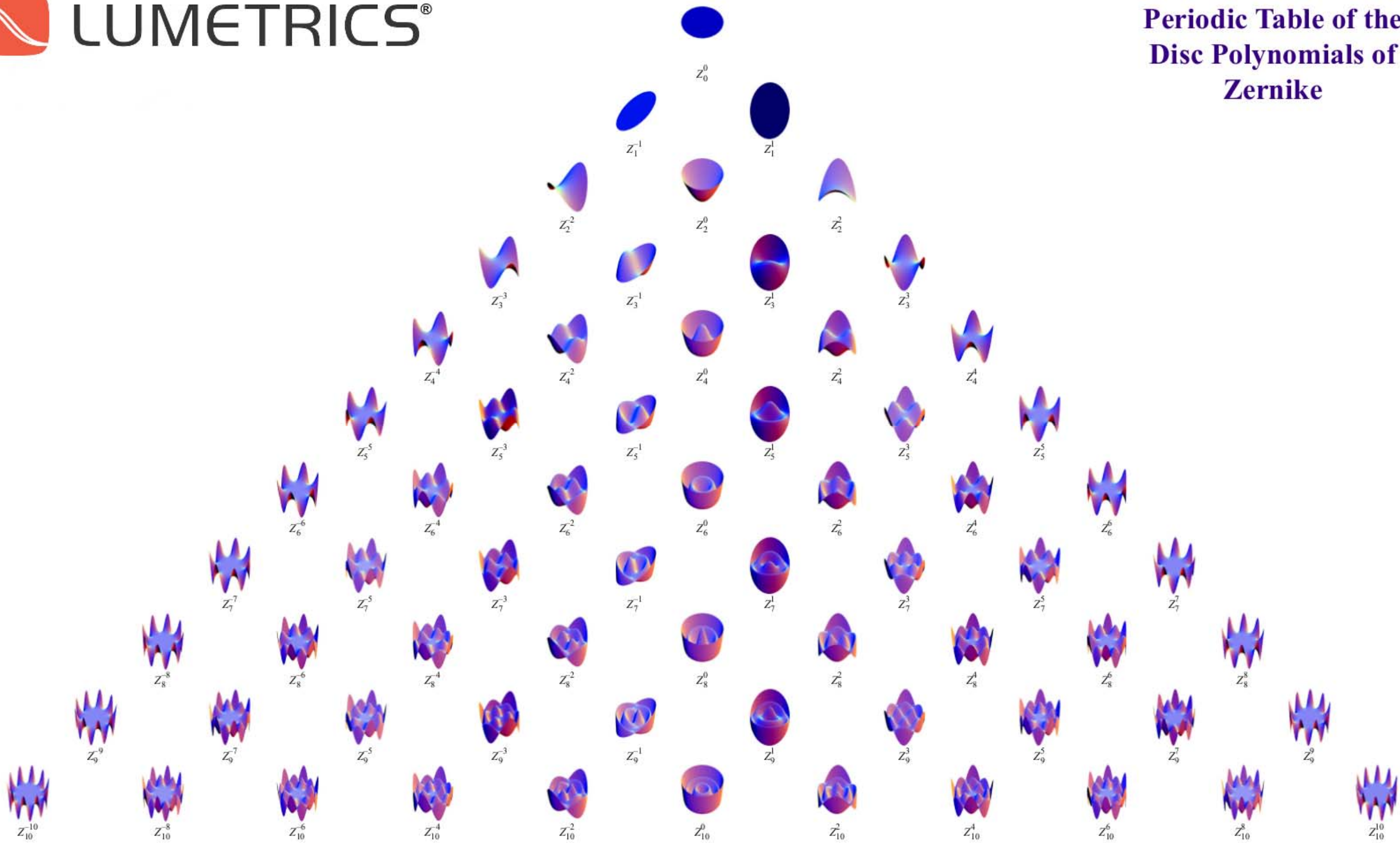
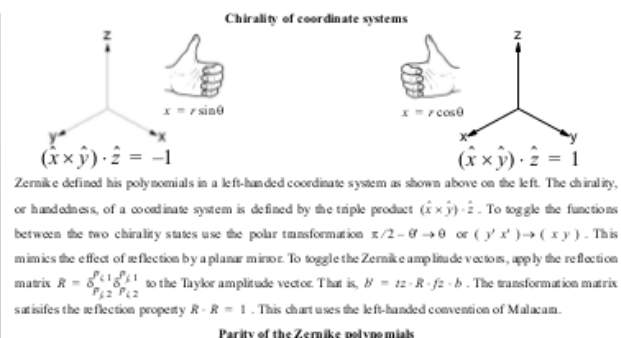


**Periodic Table of the
Disc Polynomials of
Zernike**



#	n	m	μ	Z _{n,m} ^μ	Z _{n,m}	Taylor	RMS	Polar form (dx ² → r dθ r)	Cartesian form (dx ² → dx dy)	Taylor coefficient vector
0	0	0	0	Z _{0,0} ⁰	Z _{0,0}	1	√2	1	1	(1)
1	1	-1	0	Z _{1,-1} ⁰	Z _{1,-1}	x	2	rsinθ	x	(0),(1,0)
2	1	1	1	Z _{1,1} ¹	Z _{1,1}	y	2	rcosθ	y	(0),(0,1)
3	2	-2	0	Z _{2,-2} ⁰	Z _{2,-2}	x ²	√6	r ² sin2θ	2xy	(0),(0),(2,0)
4	2	0	1	Z _{2,0} ¹	Z _{2,0}	xy	√3	2r ² -1	-1+2x ² +2y ²	((-1),(0),(0,2))
5	2	2	2	Z _{2,2} ²	Z _{2,2}	y ²	√6	r ² cos2θ	-x ² +y ²	(0),(0),(-1,0,1)
6	3	-3	0	Z _{3,-3} ⁰	Z _{3,-3}	x ³	2√2	r ³ sin3θ	-x ³ +3xy ²	(0),(0),(0),(-1,0,3,0)
7	3	-1	1	Z _{3,-1} ¹	Z _{3,-1}	x ² y	2√2	(3r ³ -2r)sinθ	-2x+3x ² +3xy ²	(0),(-2,0),(3,0,3,0)
8	3	1	2	Z _{3,1} ²	Z _{3,1}	xy ²	2√2	(3r ³ -2r)cosθ	-2y+3y ² +3xy ²	(0),(0,-2),(0),(3,0,3,0)
9	3	3	3	Z _{3,3} ³	Z _{3,3}	y ³	2√2	r ³ cos3θ	y ³ -3xy ²	(0),(0),(0),(0,-3,0,1)
10	4	-4	0	Z _{4,-4} ⁰	Z _{4,-4}	x ⁴	√10	r ⁴ sin4θ	-4x ³ +4xy ³	(0),(0),(0),(0,-4,0,4,0)
11	4	-2	1	Z _{4,-2} ¹	Z _{4,-2}	x ³ y	√10	(4r ⁴ -3r ²)sin2θ	-6xy+8x ² y+8xy ³	((0),(0),(0,-6,0),(0),(0,8,0,8,0))
12	4	0	2	Z _{4,0} ²	Z _{4,0}	x ² y ²	√5	6r ⁴ -6r ² +1	1-6x ² -6y ² +6x ⁴ +12x ² y ² +6y ⁴	((1),(0),(-6,0,-6),(0),(6,0,12,0,6))
13	4	2	3	Z _{4,2} ³	Z _{4,2}	xy ³	√10	(4r ⁴ -3r ²)cos2θ	3x ² -3y ² -4x ⁴ +4y ⁴	(0),(0),(3,0,-3),(0),(0,-4,0,0,4,0)
14	4	4	4	Z _{4,4} ⁴	Z _{4,4}	y ⁴	√10	r ⁴ cos4θ	x ⁴ -6x ² y ² +y ⁴	(0),(0),(0),(0),(1,0,-6,0,1)



Chirality of coordinate systems
 Zernike defined his polynomials in a left-handed coordinate system as shown above on the left. The chirality, or handedness, of a coordinate system is defined by the triple product $(\hat{x} \times \hat{y}) \cdot \hat{z}$. To toggle the functions between the two chirality states use the polar transformation $\pi/2 - \theta \rightarrow \theta$ or $(y' x') \rightarrow (x y)$. This mimics the effect of reflection by a planar mirror. To toggle the Zernike amplitude vectors, apply the reflection matrix $R = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$ to the Taylor amplitude vector. That is, $b' = rz \cdot R \cdot fz \cdot b$. The transformation matrix satisfies the reflection property $R \cdot R = 1$. This chart uses the left-handed convention of Malacara.

Parity of the Zernike polynomials
 The angular parity of the Zernike polynomials is determined by the sign of the angular frequency m : $Z_n^m(r, \theta) = \pm Z_n^m(r, \theta)$; when $m \geq 0$, the functions are even, when $m < 0$, the functions are odd. In lieu of resolving a wavefront into separate amplitudes for $\pm m$, consider the two terms to describe a state of mixed parity and report one amplitude and α , the angle which defines how the two definite parity states are mixed. The mixing angle is computed via $\cos \alpha = \frac{a_n^m}{\sqrt{a_n^m{}^2 + b_n^m{}^2}}$ where $k = \sqrt{a_n^m{}^2 + b_n^m{}^2}$ and $M = |m|$.

How Zernike amplitudes transform under arbitrary rotation
 The change in the Zernike amplitudes is easy to compute when the wavefront is rotated by an arbitrary angle θ . The rotation is applied to the mixed parity states, not to the definite parity states. The mixing angle is rotated using $\alpha'_m = \alpha_m + M\theta$. Clearly there is no need to rotate the rotationally invariant states $Z_{2,0}, Z_{4,0}, Z_{6,0}$.

How Zernike amplitudes transform under aperture reduction
 Computing the Zernike amplitudes for a reduced concentric aperture is trivial. In the Taylor basis, this operation is a rescaling of the coordinate axes. Consider the reduced aperture to be smaller by a multiplicative factor f ($0 < f \leq 1$). The $n+1$ amplitudes of order n are multiplied by f^n which is accomplished by multiplying the Taylor amplitudes by a $i \times i$ matrix F whose elements are given by $F_{ij} = f^{n+1} \delta_{ij}$. So given a vector b of Zernike amplitudes, the amplitudes b' for the reduced aperture are $b' = rz \cdot F \cdot fz \cdot b$.

Computing wavefront error using the Zernike amplitudes
 Given a vector of Zernike amplitudes b representing a j th order expansion, one can compute σ^2 the wavefront error (or wavefront variance) using $\sigma^2 = \frac{1}{\pi} \int_0^{2\pi} \int_0^1 |b \cdot Z_n(r, \theta)|^2 r dr d\theta$.

There is no contribution from the piston term since this is the mean value \bar{w} . For example, in a second order expansion $\sigma^2 = \frac{1}{4}(b_{1,0}^2 + b_{1,1}^2 + \frac{1}{2}b_{2,0}^2 + 2b_{2,1}^2 + b_{2,2}^2)$. The polynomials can be scaled by the factors in the RMS column to simplify the wavefront variance to $\sigma^2 = b_{1,0}^2 + b_{1,1}^2 + b_{2,0}^2 + b_{2,1}^2 + b_{2,2}^2$.

The radial polynomials of Zernike: the foundation for the disk polynomials
 Symbol: $R_n^m(r)$ or $R_{n,\mu}(r)$; where $m = 2\mu - n$ and $\mu = \frac{n+m}{2}$. E.g., $R_4^0(r) = R_{4,2}(r) = 6r^4 - 6r^2 + 1$
 Interval: $[0, 1]$ Weight: 1 Inequality: $|R(r)| \leq 1$ Standardization: $R(1) = 1$ Norm: $1/(2n+2)$

Differential equation: $r(1-r^2)R'' + (1-3r^2)R' + (n(n+2)r - \frac{m^2}{r})R = 0$ where $R = R_n^m(r)$
 Explicit expression: $R_n^m(r) = \sum_{l=0}^n \frac{(-1)^l (n-l)!}{l! [\frac{1}{2}(n+m-l)]! [\frac{1}{2}(n-m-l)]!} r^{n-l}$
 Recurrence relations: $R_n^m(r) = \frac{1}{2(n+1)r} [(n+m+2)R_{n+1}^{m+1}(r) + (n-m)R_{n+1}^{m-1}(r)]$
 $R_{n+2}^m(r) = \frac{n+2}{(n+2)^2 - m^2} [4(n+1)r^2 - \frac{(n+m)^2}{n} - \frac{(n-m+2)^2}{n+2}] R_n^m(r) - \frac{n^2 - m^2}{n} R_{n-2}^m(r)$
 $R_n^m(r) + R_{n+2}^m(r) = \frac{1}{n+1} \frac{d[R_{n+1}^{m+1}(r) - R_{n+1}^{m-1}(r)]}{dr}$

Relationships to other special functions:
 Jacobi polynomial: $R_n^m(r) = (-1)^{\frac{n-m}{2}} r^{\frac{m}{2}} F_4^{(m,0)}(1-r^2)$ where $P_n^{(\alpha,\beta)}(x) = \frac{1}{2} \sum_{k=0}^n \binom{n+\alpha}{k} \binom{n+\beta}{n-k} (x-1)^k (x+1)^{n-k}$
 Hypergeometric functions: $R_n^m(r) = (-1)^{\frac{n-m}{2}} \frac{1}{2} {}_2F_1(\frac{n+m+2}{2}, \frac{n-m}{2}, m+1, r^2)$ where ${}_2F_1(a, b; c; x) = \sum_{n=0}^{\infty} \frac{(a)_n (b)_n}{(c)_n} \frac{x^n}{n!}$

Bessel function of the first kind: $\int_0^1 R_n^m(r) J_m(xr) r dr = (-1)^{\frac{n-m}{2}} \frac{J_{n+1}(x)}{x}$ where $J_n(x) = \sum_{l=0}^{\infty} \frac{(-1)^l}{2^{2l} l! (n-l)!} x^{2l-n}$
Legendre polynomials: $P_n(r) = 2^n \sum_{l=0}^n (-1)^l \binom{n}{l} \binom{2n-2l}{n} r^{n-2l}$ resemble $R_n^0(r) = \sum_{l=0}^n (-1)^l \binom{n}{l} \binom{n-2l}{n} r^{n-2l}$ where n is even

These special functions can be generated by performing a Gram-Schmidt orthogonalization of a power series

Functions	Series	Interval	Weight	Norm
Legendre	$\{1, r, r^2, r^3, \dots\}$	$-1 \leq r \leq 1$	1	$2/(2n+1)$
Shifted Legendre	"	$0 \leq r \leq 1$	1	$1/(2n+1)$
Chebyshev I	"	$-1 \leq r \leq 1$	$(1-x^2)^{-1/2}$	$\pi/(2-\delta_0^2)$
Shifted Chebyshev I	"	$0 \leq r \leq 1$	$[x(1-x^2)]^{-1/2}$	$\pi/(2-\delta_0^2)$
Chebyshev II	"	$-1 \leq r \leq 1$	$(1-x^2)^{1/2}$	$\pi/2$
Associated Laguerre	"	$0 \leq r < \infty$	$r^k e^{-r}$	$(n+k)!/n!$
Hermite	"	$-\infty < r < \infty$	e^{-r^2}	$2^n \sqrt{\pi} n!$
Zernike radial	$\{r^m, r^{m+2}, r^{m+4}, \dots\}$	$0 \leq r \leq 1$	r	$1/(2n+2)$

Important properties of the circle polynomials of Zernike

- Completeness:** $\lim_{D_2 \rightarrow \infty} \int_{D_2} \psi(x_1, x_2) \sum_{n=0}^{\infty} \sum_{m=-n}^n b_{n,m} Z_{n,m}(x_1, x_2) dx^2 = 0$
 The polynomial approximation converges uniformly to the input function (Weierstrass theorem).
- Orthogonality:** $\int_{D_2} Z_{n,m}(x_1, x_2) Z_{n',m'}(x_1, x_2) dx^2 = \frac{1 + \delta_{n,n'} \pi \delta_{m,m'}}{2(n+1)}$. The amplitudes do not change as the order of the fit is increased.
- The Cartesian Zernike polynomials are exactly equivalent to a Taylor series order by order.**
 Given a set of Zernike amplitudes $b_{n,m}$ there is an exact affine transformation to and from the Taylor amplitudes $a_{n,m}$. If fz is the transformation from the Zernike basis to the Taylor basis and rz is the transformation from the Taylor basis to the Zernike basis,

$$a = fz \cdot b, b = rz \cdot a. \text{ E.g., the transformation matrices for second order are } fz = \begin{bmatrix} 1 & 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 2 & 1 \end{bmatrix}, rz = \begin{bmatrix} 1 & 0 & 0 & 1/4 & 0 & 1/4 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 2 \\ 0 & 0 & 0 & 1/4 & 0 & 1/4 \\ 0 & 0 & 0 & -1/2 & 0 & 1/2 \end{bmatrix}$$

The Zernike functions are replaced by vectors in the Taylor basis. The Taylor series $\{1, x, y, x^2, xy, y^2, \dots\}$ generalizes to $x^i y^j r^k$ with the vector of exponential powers $p = (0, 0, (1, 0), (0, 1), (2, 0), (1, 1), (0, 2), \dots)$.

- Addressing and indexing**
- If the collection of Zernike coefficient vectors is z then $z[n]$ is the coefficient vectors for order n ; $z[n][m]$ is the coefficient vectors for the m th function of order n ; $z[\mu][l]$ is the coefficient for the l th term in the μ th function of order n .
 - A coefficient vector for order n has $z = (n+1)(n+2)/2$ terms; a coefficient vector with z terms represents order $n = \frac{1}{2}(\sqrt{4z+3}-3)$.
 - The index number of the amplitude for $Z_{n,m}(x_1, x_2)$ is $j = n(n+1)/2 + \mu + 1$; the inverse problem is solved by a lookup vector $q = ((0, 0), (1, 0), (1, 1), (2, 0), (2, 1), (2, 2), \dots)$ which provides the order n and cardinal number m .
 - The order is given by n ; the frequency is given by m ; the ordinal number is μ .

Example: Primary aberrations through a pupil of radius r :
 Taylor basis: $sphere = \frac{a_{2,0} + a_{2,2}}{2} + cyl, cyl = \frac{2}{\sqrt{2}} \sqrt{a_{2,0}^2 + a_{2,1}^2 - 2a_{2,0}a_{2,1} + a_{2,2}^2}, axis = \frac{1}{2} \arctan\left(\frac{a_{2,1}}{a_{2,2} - a_{2,0}}\right)$
 Zernike basis: $sphere = \frac{4}{3} b_{2,0} + cyl, cyl = \frac{4}{\sqrt{2}} \sqrt{b_{2,0}^2 + b_{2,1}^2}, axis = \frac{1}{2} \arctan\left(\frac{b_{2,1}}{b_{2,2}}\right)$

The circle polynomials of Zernike are defined over the unit disk D_2 . In polar coordinates $0 \leq r \leq 1, 0 \leq \theta < 2\pi$; in Cartesian coordinates $x^2 + y^2 \leq 1$

Connection between the radial polynomials $R_n^m(r)$ and the circle polynomials $Z_n^m(r, \theta)$: The complex combination $Z_n^m(r, \theta) \pm i Z_n^m(r, \theta) = r^n e^{im\theta} = R_n^m(r) e^{im\theta}$ leads to $Z_n^m(r, \theta) = R_n^m(r) \cos m\theta$ and $Z_n^m(r, \theta) = R_n^m(r) \sin m\theta$
Polar recursion relationship: $R_n^m(r) = \sum_{j=0}^n (-1)^j \frac{(n-j)!}{j!(n-m-j)!(m-j)!} r^{n-2j}$, Cartesian recursion: $Z_n^m(x, y) = \sum_{i=0}^n \sum_{j=0}^{M-i} (-1)^j \frac{(n-2M)!}{(2i+p)!} \frac{(M-j)!}{k!(M-j)!(n-M-j)!} x^i y^j$ where $n = n-2m, s = \text{sgn}(d), M = \frac{n}{2} - |m - \frac{n}{2}|, p = \frac{|j|}{2}(s+1), q = (d - i \text{Mod}[n, 2]) \frac{1}{2}, \xi = 2(i+k) + p, j = n - 2(i+j+k) - p$